



Charge Symmetry Breaking in the $dd \rightarrow {}^4\text{He}\pi^0$ Reaction with WASA-at-COSY

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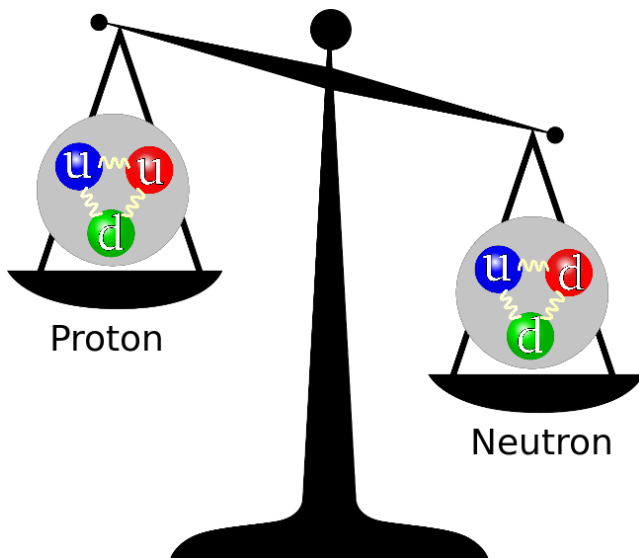
Two sources of violation:

- Electromagnetic interaction
- Lightest quark mass difference \mapsto Window for probing quark mass ratios

Isospin Symmetry

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Nucleon mass difference

$$\Delta M_{np} = \Delta M_{em} + \Delta M_{str}$$

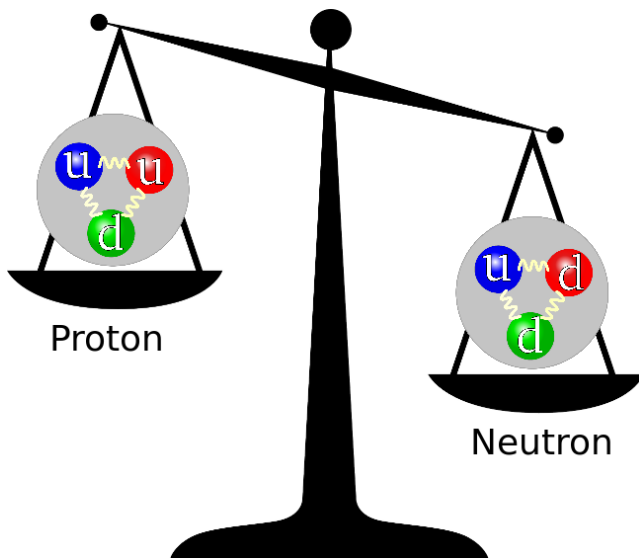
$-0.7 \pm 0.3 \text{ MeV}$
 (from QED + dispersion theory)

$2.05 \pm 0.3 \text{ MeV}$
 $(\Delta M_{pn} - \Delta M_{em})$

Isospin Symmetry

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- Electromagnetic interaction
- Lightest quark mass difference \rightarrow Window for probing quark mass ratios



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$-0.7 \pm 0.3 \text{ MeV}$
(from QED + dispersion theory)

$2.05 \pm 0.3 \text{ MeV}$
($\Delta M_{pn} - \Delta M_{em}$)

Access to ΔM_{str} from dynamic ISB from Chiral Perturbation Theory

πN scattering length, e.g., $a(\pi^0 p) - a(\pi^0 n) = f(\Delta M_{str})$ (Weinberg 1977)

However:

- No direct measurement of $\pi^0 N$
- Large e.m. corrections in $\pi^\pm N$

Charge Symmetry Breaking

Isospin Symmetry Breaking

Dominated by pion mass difference Δm_π – e.m. effect



Charge Symmetry (CS) Breaking

Symmetry under the operation of $P_{CS} = e^{-i\tau_2\pi/2}$ - Δm_π does not contribute

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Charge Symmetry (CS) Breaking

Symmetry under the operation of $P_{CS} = e^{-i\tau_2\pi/2}$ – Δm_π does not contribute

1. $np \rightarrow d\pi^0$ forward-backward asymmetry A_{fb} [1]

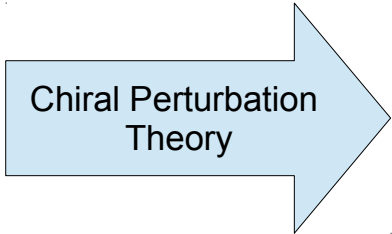
$$\Delta M_{str} = (1.5 \pm 0.8 \text{ (exp.)} \pm 0.5 \text{ (th.)}) \text{ MeV (LO) [2]}$$

2. $dd \rightarrow {}^4\text{He}\pi^0$

$$\text{CS} \Rightarrow \sigma = 0 \quad \text{CSB} \Rightarrow \sigma \neq 0, \sigma \propto |M_{CSB}|^2 = |M_1 + M_2 + \dots|^2$$

σ_{total} measured at threshold [3] and at $Q = 60 \text{ MeV}$ [4]

Result at threshold
consistent with s-wave



Chiral Perturbation
Theory

p-wave contribution in $dd \rightarrow {}^4\text{He}\pi^0$
at higher excess energies needed

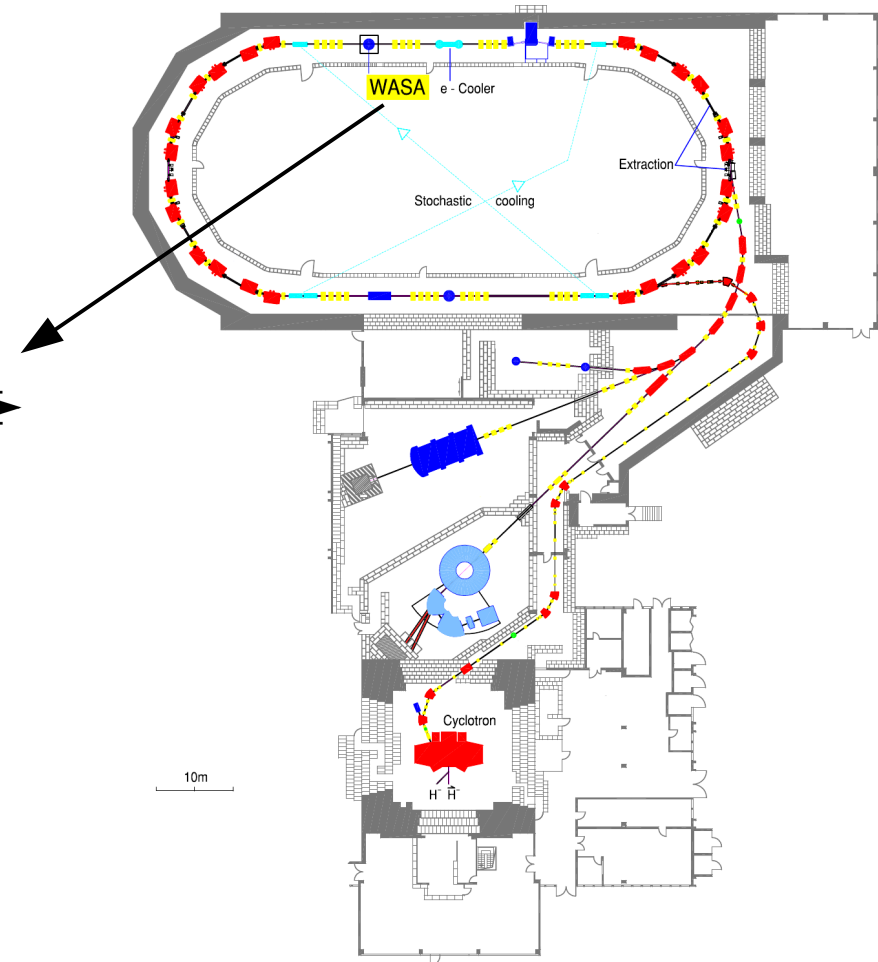
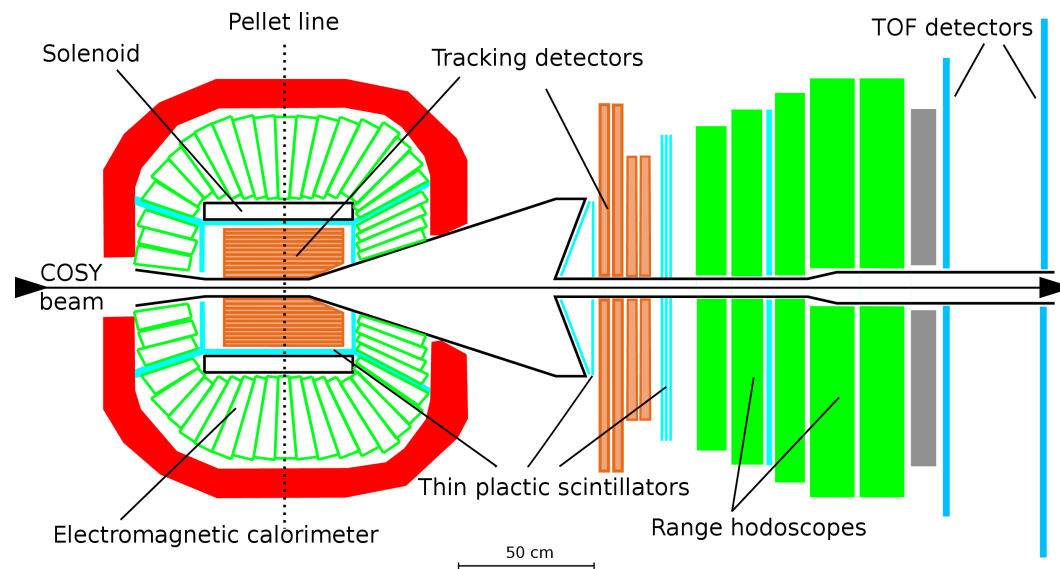
[1] Opper et al. PRL 91 (2003) 212302

[3] Stephenson et al. PRL 91 (2003) 142302

[2] Filin et al. Phys. Lett. B681 (2009) 423

[4] Adlarson et al. Phys. Lett. B 739 (2014) 44

WASA-at-COSY experiment



CSB with WASA-at-COSY:

2007: Measurement of $dd \rightarrow {}^3\text{He}n\pi^0$

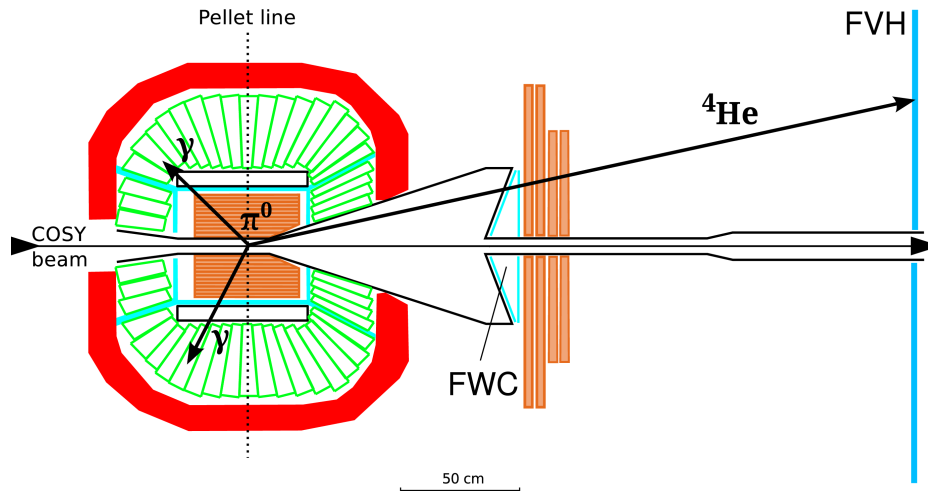
goal: description of main background, input for initial-state-interaction calculations

2008: First measurement of $dd \rightarrow {}^4\text{He}n\pi^0$ (2 weeks) @ $Q = 60$ MeV

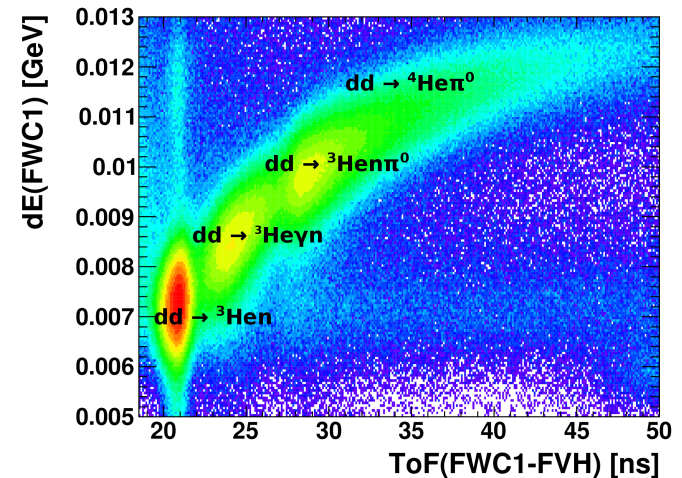
goal: σ_{total}

2014: New measurement of $dd \rightarrow {}^4\text{He}n\pi^0$ (8 weeks) @ $Q = 60$ MeV with modified detector
goal: angular distribution

Analysis of $dd \rightarrow {}^4\text{He}\pi^0$



Status after calibration:



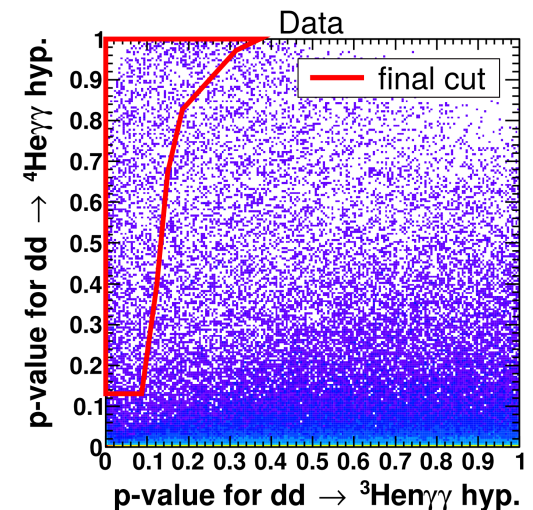
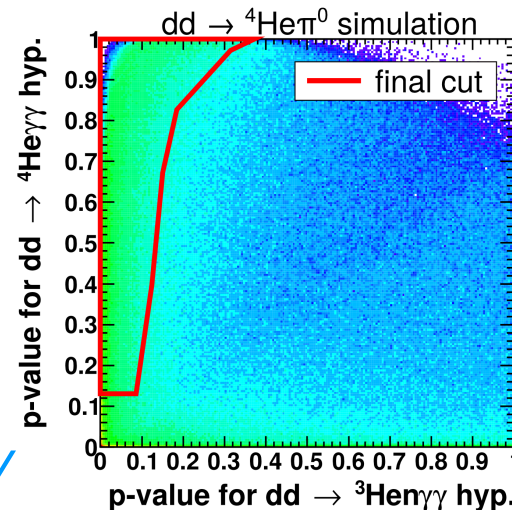
Background

- $dd \rightarrow (pnd, pnpn, tp) + \pi^0$
- $dd \rightarrow {}^3\text{He}n\pi^0$ ($3 \cdot 10^5$ higher σ)
- $dd \rightarrow {}^4\text{He}\gamma\gamma$ (physics bg)

Overall kinematic fit

→ 2 hypotheses fitted:

$dd \rightarrow {}^4\text{He}\gamma\gamma$ and $dd \rightarrow {}^3\text{He}n\gamma\gamma$



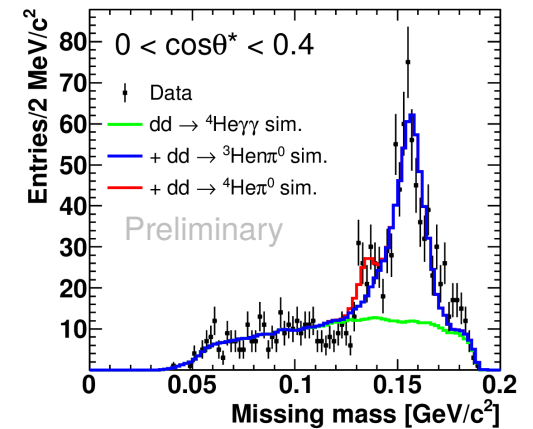
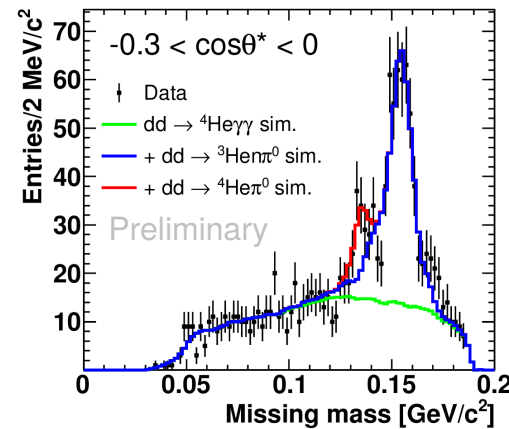
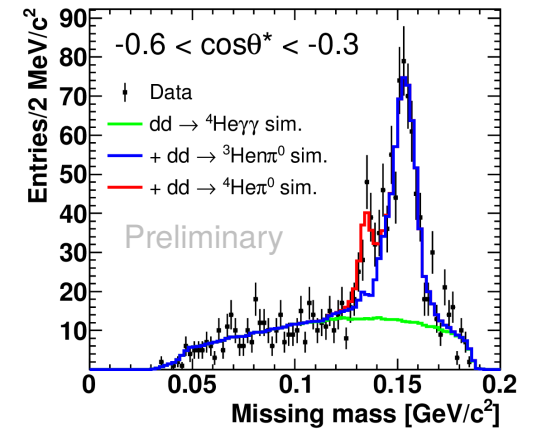
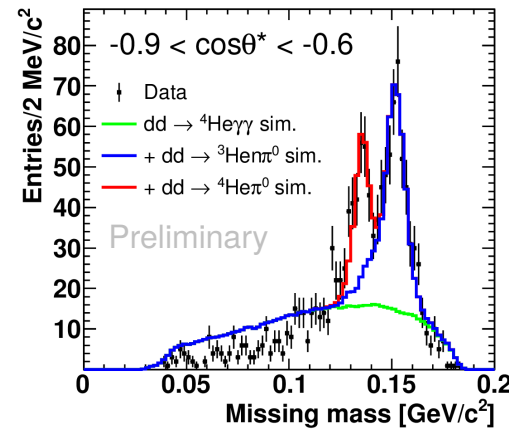
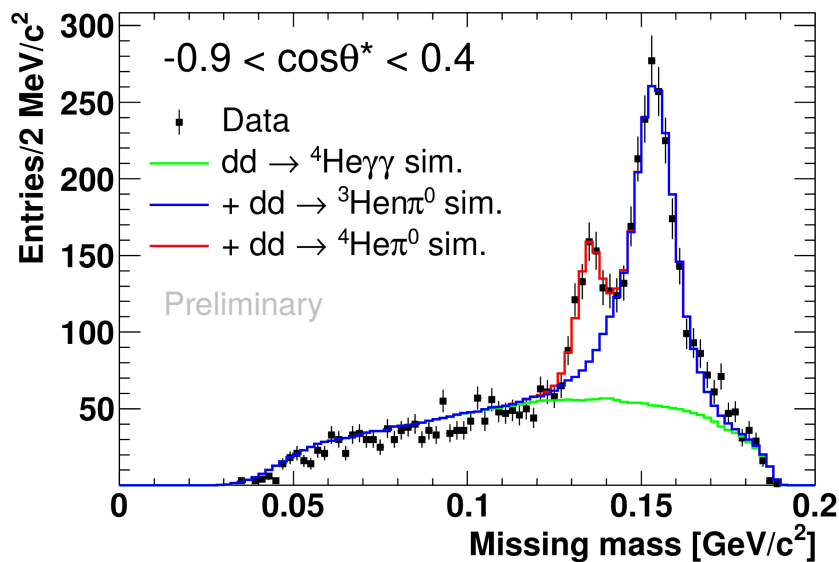
→ Optimized cuts on cumulated probability distribution (p-value)

→ Suppression of $dd \rightarrow {}^3\text{He}n\pi^0$ about 10^4

Missing mass of $dd \rightarrow {}^4\text{He}X$

Four angular bins

Full angular range
within detector acceptance

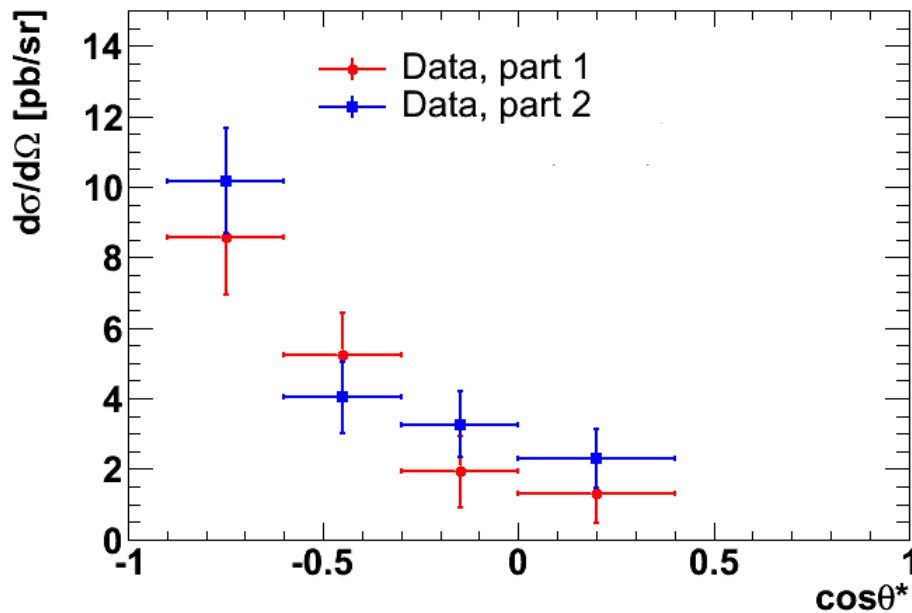


Luminosity determination using $dd \rightarrow {}^3\text{He}\pi^0$

Total and differential cross section

Unpolarized differential cross section (terms up to order $p_{\pi^0}^2$ to the intensity):

$$\frac{d\sigma}{d\Omega} = \frac{2 p_{\pi^0}}{3 p_d} \left(|A_0|^2 - p_{\pi^0}^2 \Re\{A_0^* A_2\} + |C|^2 p_{\pi^0}^2 \right) + \frac{p_{\pi^0}}{p_d} \left(2 p_{\pi^0}^2 \Re\{A_0^* A_2\} - \frac{2}{3} |C|^2 p_{\pi^0}^2 \right) \cos^2 \theta^*$$

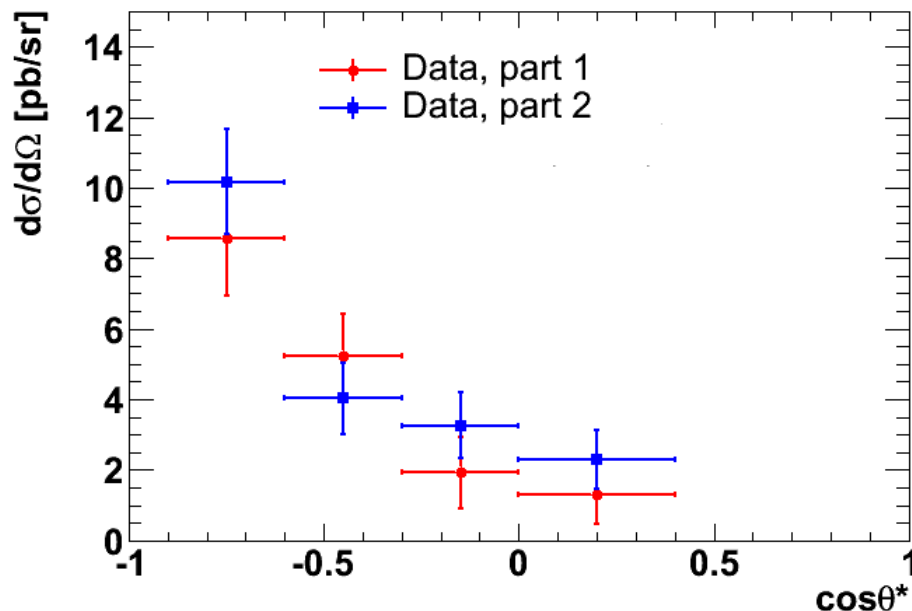


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s-wave amplitude



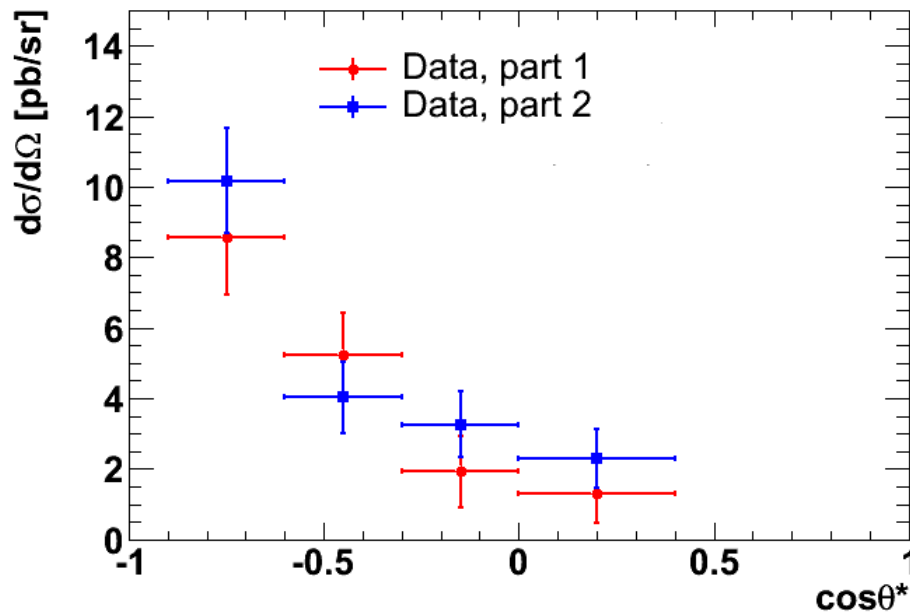
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s-wave amplitude

s-*d* interference term



Total and differential cross section

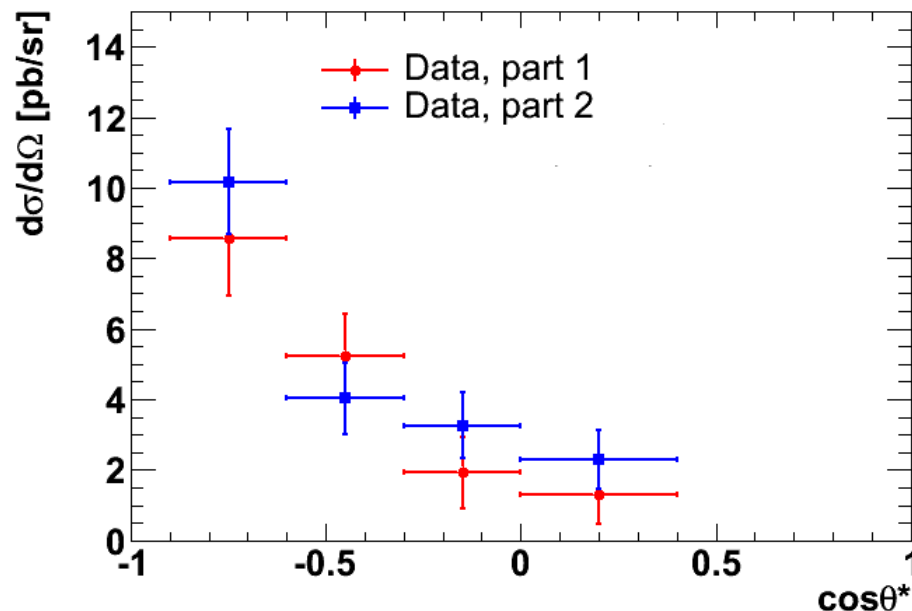
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s-wave amplitude

p-wave amplitude

s-d interference term



Total and differential cross section

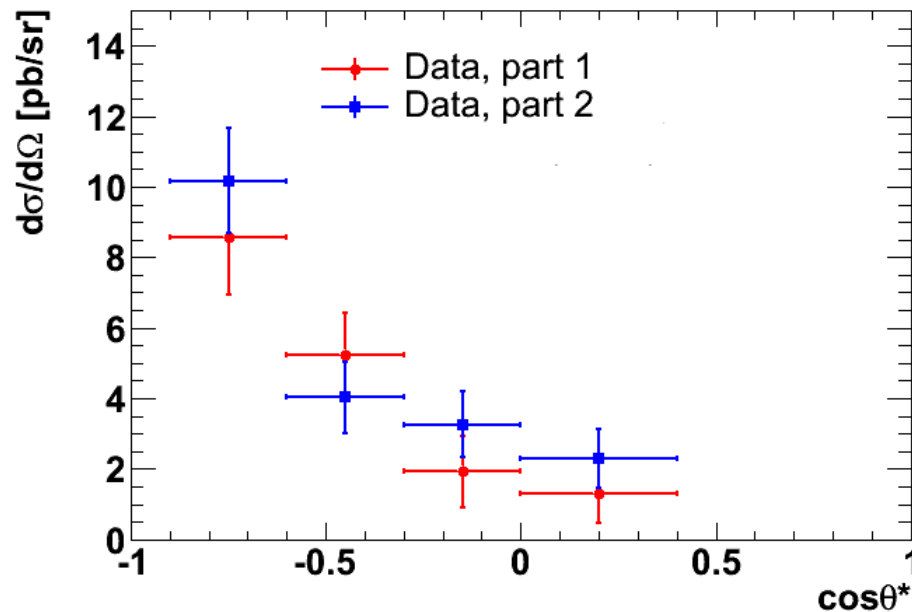
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s-wave amplitude

p-wave amplitude

s-d interference term



Total and differential cross section

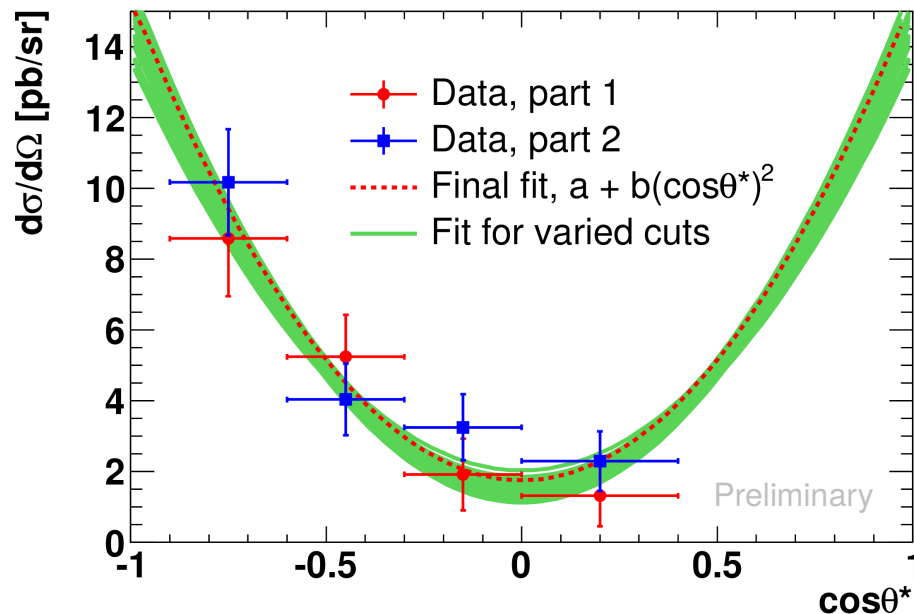
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s-wave amplitude

p-wave amplitude

s-d interference term



$$a = \left(1.75 \pm 0.46(\text{stat.})_{-0.8}^{+0.31}(\text{syst.}) \right) \text{ pb/sr}$$

$$b = \left(13.6 \pm 2.2(\text{stat.})_{-2.7}^{+0.9}(\text{syst.}) \right) \text{ pb/sr}$$

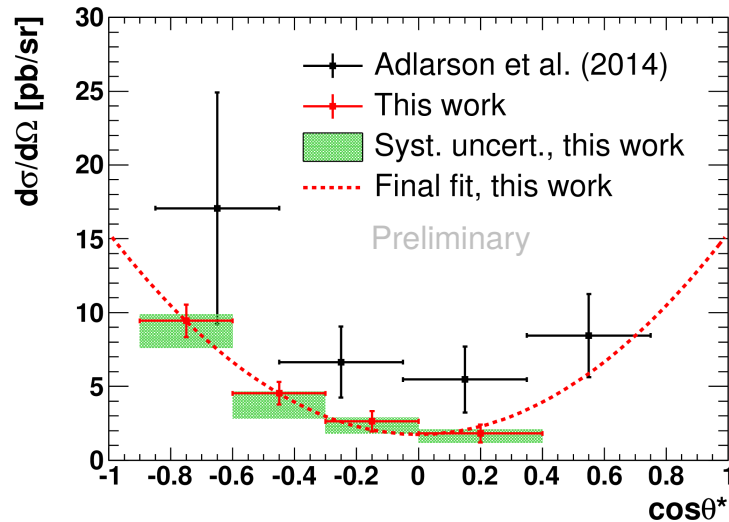
common systematic uncertainties

10% - external normalization

2% - luminosity determination

→ Requires higher partial waves

$$\sigma_{\text{tot}} = \left(79.1 \pm 7.3(\text{stat.})_{-10.5}^{+1.2}(\text{syst.}) \pm 8.1(\text{norm.}) \pm 2.0(\text{lumi. syst.}) \right) \text{ pb}$$



$$\sigma_{\text{tot}} = \left(79.1 \pm 7.3(\text{stat.})^{+1.2}_{-10.5}(\text{syst.}) \pm 8.1(\text{norm.}) \pm 2.0(\text{lumi. syst.}) \right) \text{ pb}$$

$$\sigma_{\text{tot}}^{\text{prev}} = (123 \pm 30(\text{stat.}) \pm 12(\text{norm.}) \pm 8.6(\text{ext.})) \text{ pb}$$

From energy dependence of total cross section we can obtain s-wave amplitude $|A_0|^2$

(Neglecting initial and final state interactions)

$$\frac{p}{p_{\pi^0}} \sigma_{\text{tot}} = \frac{8\pi}{3} |A_0|^2 + \frac{16\pi}{9} m_{\pi^0}^2 |C|^2 \eta^2$$

Defined mostly by measurement close to threshold

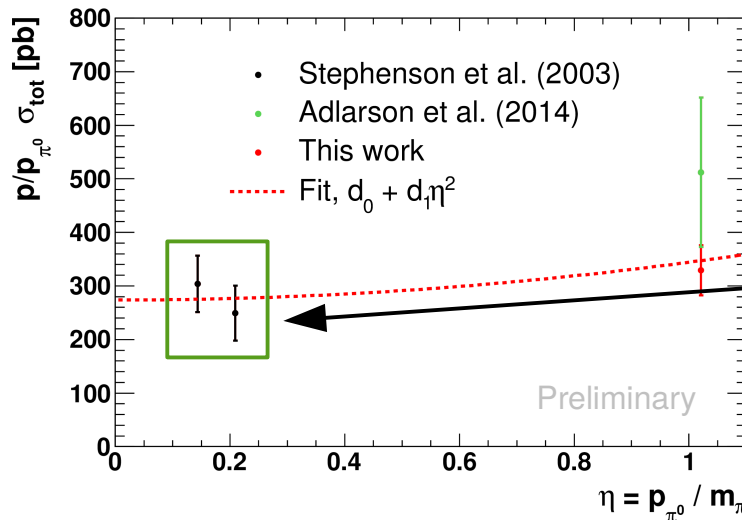
$$|A_0|^2 = (32.7 \pm 4.5) \text{ pb/sr}$$

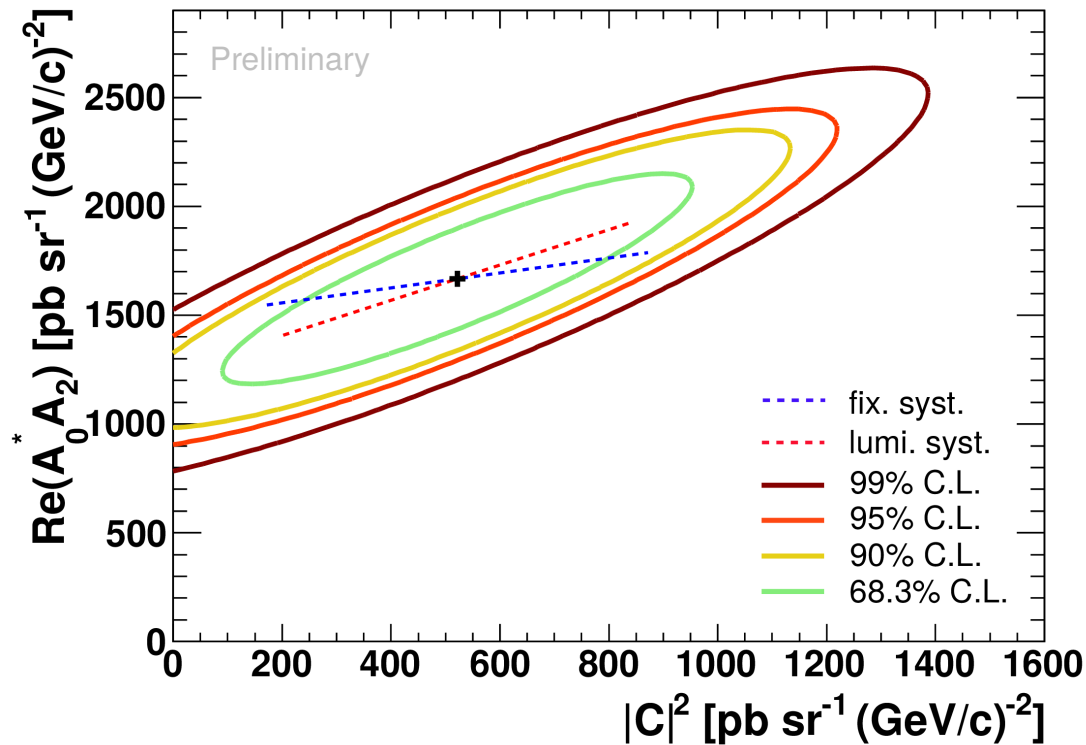
With fixed $|A_0|^2$ we can determine $|C|^2$ and $\Re\{A_0^* A_2\}$

from the fit:

$$\frac{p_d}{p_{\pi^0}} \frac{d\sigma}{d\Omega} = \frac{2}{3} \left(\boxed{|A_0|^2} - p_{\pi^0}^2 \boxed{\Re\{A_0^* A_2\}} + \boxed{|C|^2} p_{\pi^0}^2 \right) + \left(2p_{\pi^0}^2 \boxed{\Re\{A_0^* A_2\}} - \frac{2}{3} \boxed{|C|^2} p_{\pi^0}^2 \right) \cos^2 \theta^*$$

fixed





***p*-wave amplitude**

small and consistent with zero
within uncertainties

***s*-*d* interference term**

significantly different from zero

Direct contribution from $|A_2|^2$

Subtracted in σ_{tot} by factor $1/p_{\pi^0}^2 \approx 50$
(it is of order of $p_{\pi^0}^4$)

$$\Re\{A_0^* A_2\} = (1670 \pm 320(\text{stat.})_{-430}^{+80}(\text{syst.})) \text{ pb/}(\text{sr} \cdot (\text{GeV}/c)^2)$$

$$|C|^2 = (520 \pm 290(\text{stat.})_{-430}^{+50}(\text{syst.})) \text{ pb/}(\text{sr} \cdot (\text{GeV}/c)^2)$$

Common (correlated) systematic uncertainties:

- From luminosity determination
- From fixed $|A_0|^2$

Summary

- **Charge Symmetry Breaking** used to access quark mass effects.
Theoretical tool: Chiral Perturbation Theory.
- Higher partial wave contributions in $dd \rightarrow {}^4\text{He}\pi^0$ needed.
- Data from the new measurement of $dd \rightarrow {}^4\text{He}\pi^0$ at $Q = 60$ MeV in 2014 with WASA analyzed. Total and differential cross section obtained.
- **Results show that any theoretical attempt to describe the reaction had to include, in addition to p -waves, also d -wave contributions.**

Backup

Charge Symmetry Breaking

Measurements of CSB observables

- $np \rightarrow d\pi^0$ forward-backward asymmetry A_{fb}
 - leading CSB term: πN rescattering
 - Opper et al., $A_{fb} = (17.2 \pm 8.0 \pm 5.5) \cdot 10^{-3}$
(PRL 91 (2003) 212302)

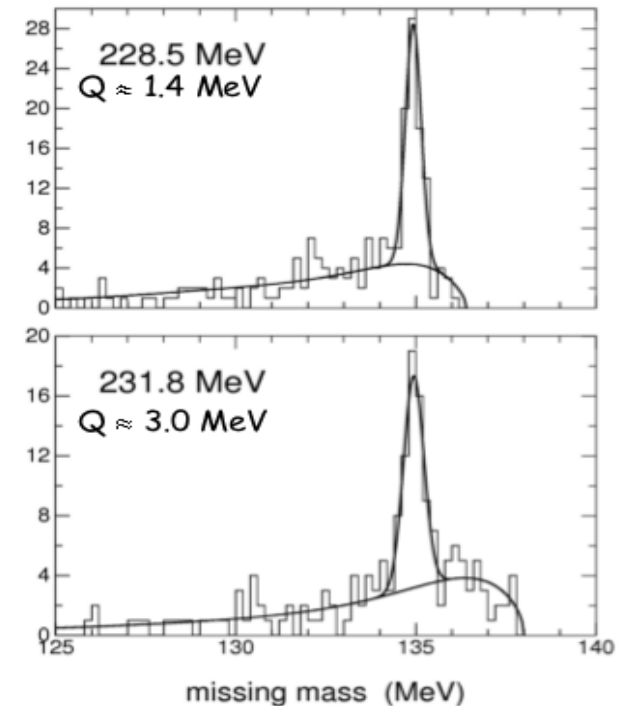
- Pion production in $dd \rightarrow {}^4\text{He} \pi^0$

CSC $\Rightarrow \sigma = 0$

CSB $\Rightarrow \sigma \neq 0, \sigma \propto |M_{CSB}|^2$

Complementary to $np \rightarrow d\pi^0$:

- different strength of CSB terms
- dd initial state more demanding



Result: Stephenson et al.

(PRL 91 (142302) 2003)

$\sigma_{tot} (Q=1.4 \text{ MeV}) = 12.7 \pm 2.2 \text{ pb}$

$\sigma_{tot} (Q=3.0 \text{ MeV}) = 15.1 \pm 3.1 \text{ pb}$

Result consistent with s-wave production

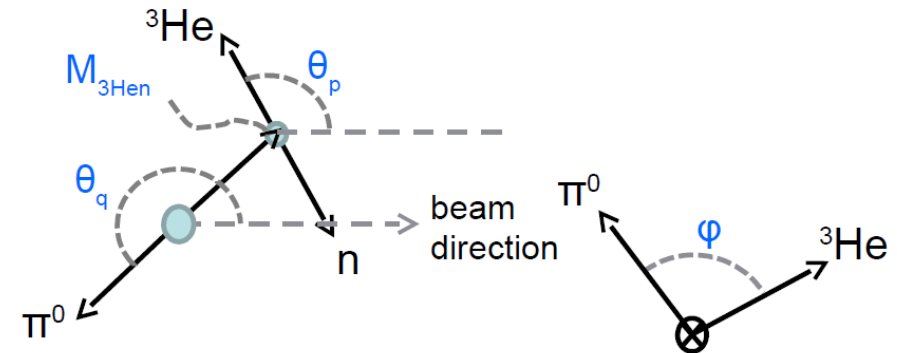
$dd \rightarrow {}^3\text{He}n\pi^0$ reaction measurement

Two-fold model ansatz:

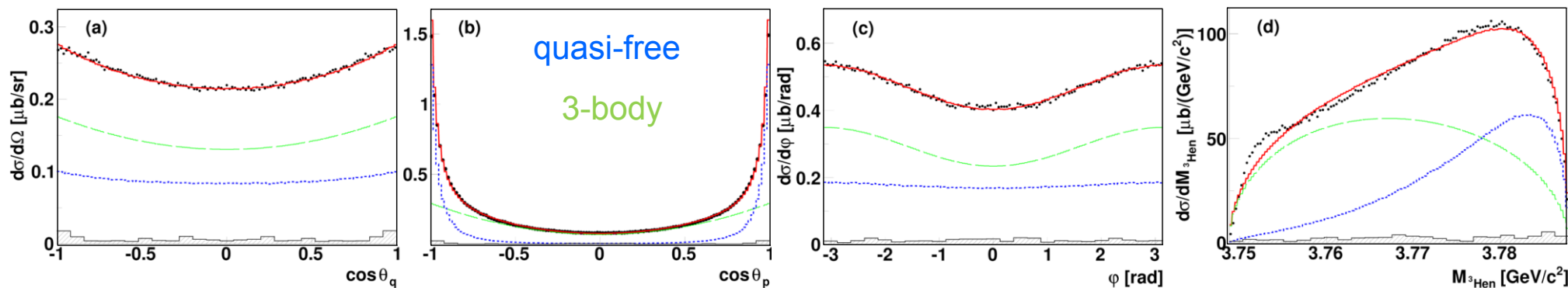
- Quasi-free contribution: $dd \rightarrow {}^3\text{He}n\pi^0 + n_{\text{spec}}$
 - Partial waves decomposition of the 3-body final state (limited to $L \leq 1$)
- full model
incoherent sum

$$\sigma_{\text{tot}} = (2.89 \pm 0.01_{\text{stat}} \pm 0.06_{\text{sys}} \pm 0.29_{\text{norm}}) \mu\text{b}$$

Model used for **simulating** the $dd \rightarrow {}^3\text{He}n\pi^0$ background and for **normalization**

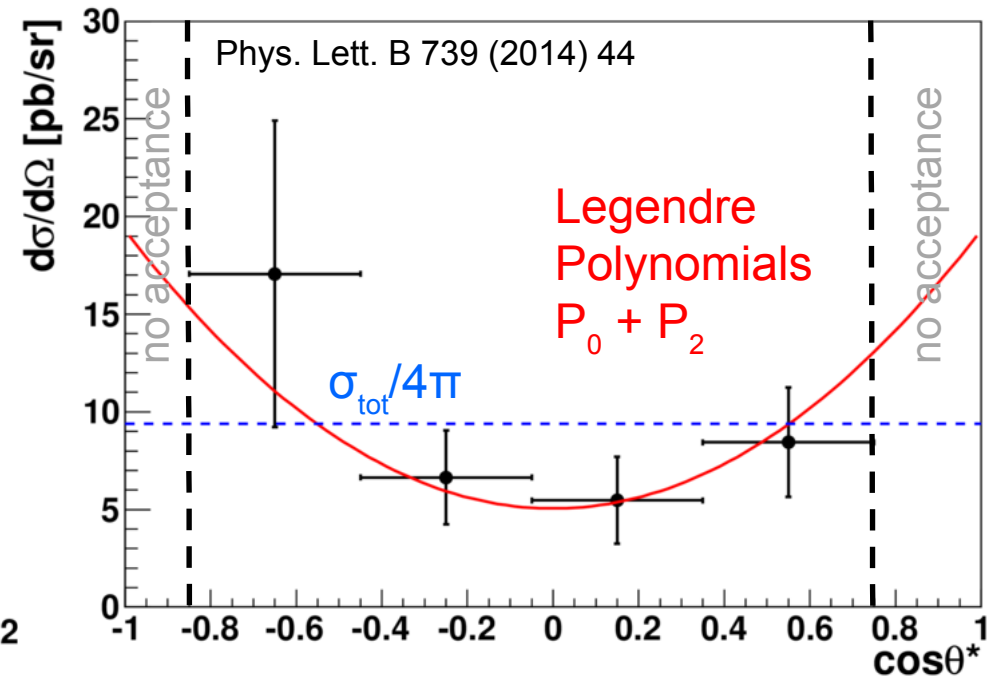
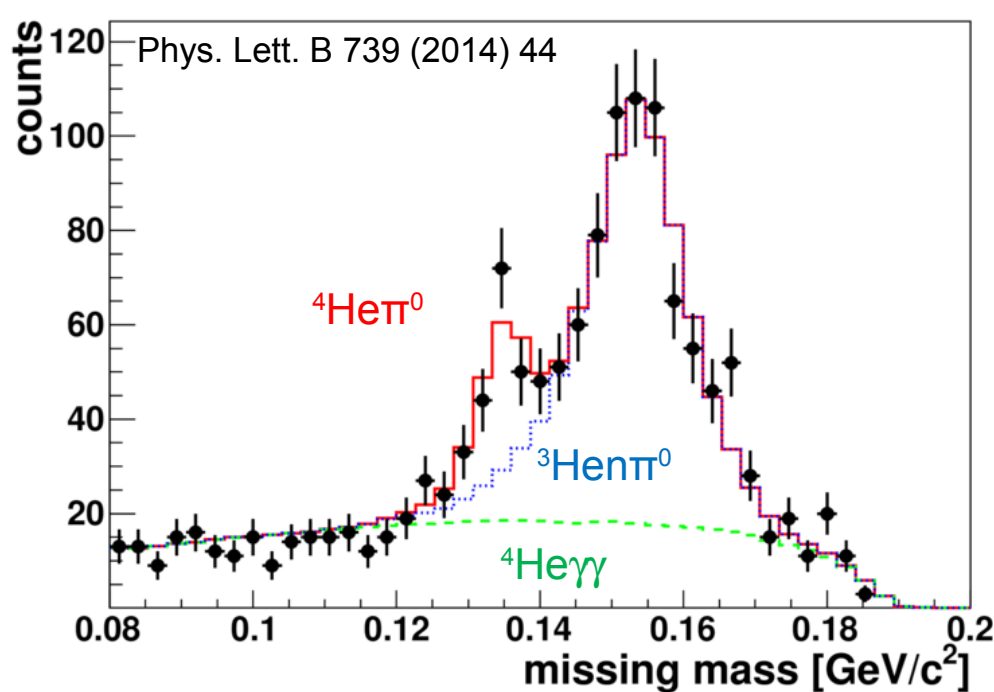


4 independent variables $M_{3\text{He}n}, \theta_p, \theta_q, \phi$



Phys. Rev. C 88 (2013) 014004

Results:



Total cross section:

$${}^4\text{He}\pi^0: \sigma = (118 \pm 18_{\text{stat}} \pm 13_{\text{sys}} \pm 8_{\text{ext}}) \text{ pb}$$

$${}^4\text{He}\gamma\gamma: \sigma = (920 \pm 70_{\text{stat}} \pm 100_{\text{sys}} \pm 70_{\text{ext}}) \text{ pb}$$

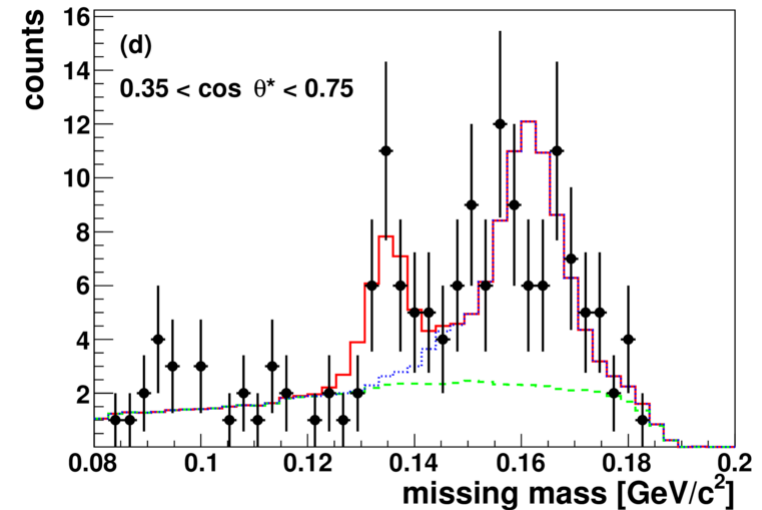
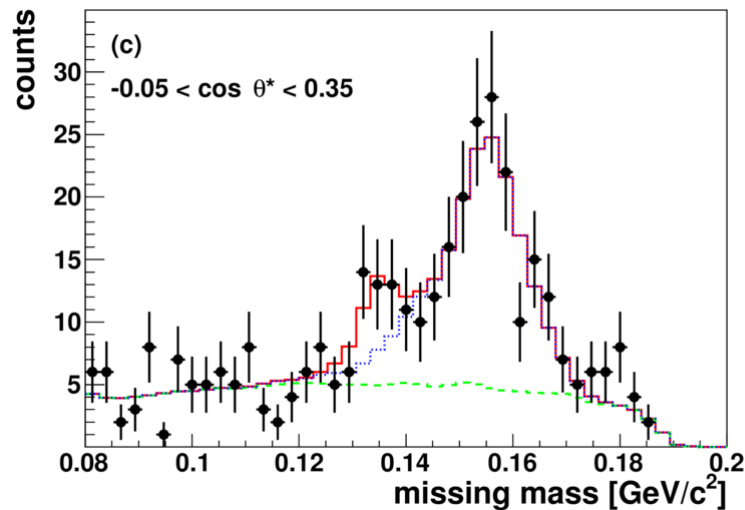
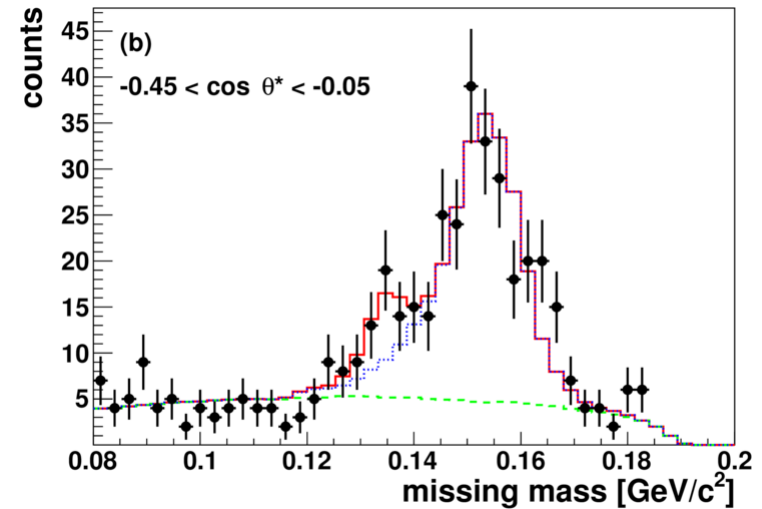
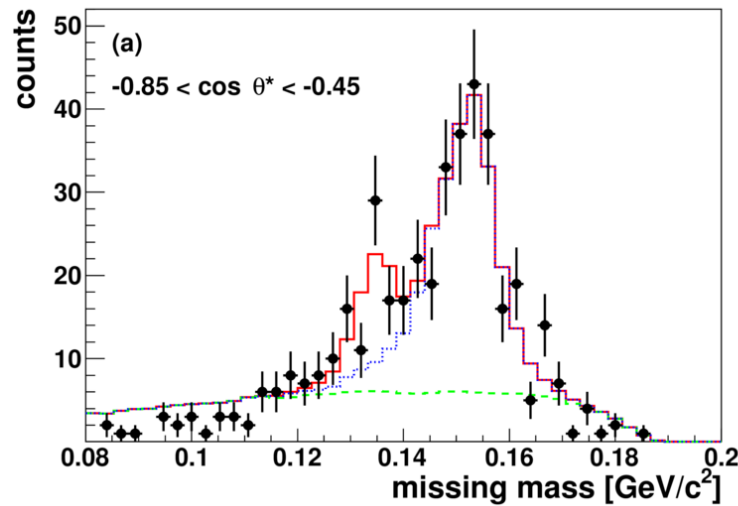
Fit including p -wave:

$$\begin{aligned} d\sigma/d\Omega = & (9.8 \pm 2.6) \text{ pb/sr} \cdot P_0(\cos\theta^*) \\ & + (9.5 \pm 7.4) \text{ pb/sr} \cdot P_2(\cos\theta^*) \end{aligned}$$

consistent with s -wave only

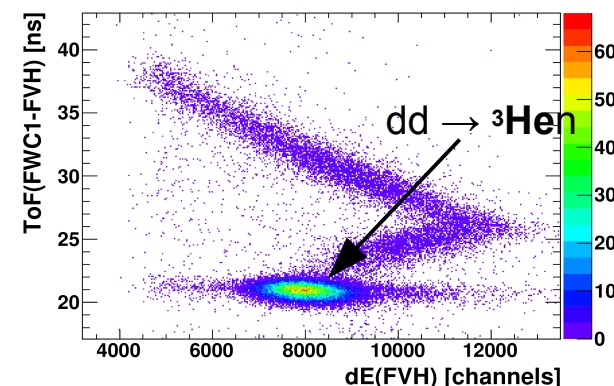
However: not decisive due to limited statistics

Results – angular distribution



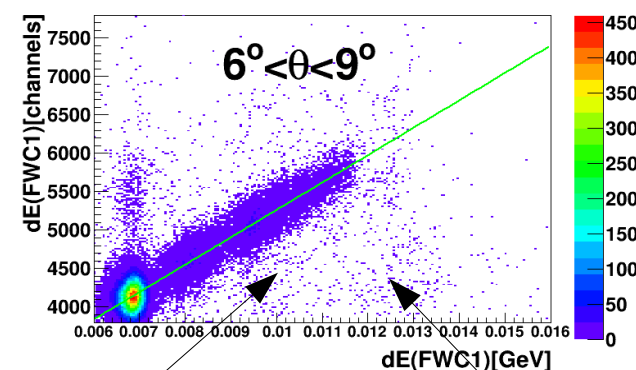
ToF Calibration

- $dd \rightarrow {}^3\text{He}n$ time peak position used
- Calibrate the data to the MC values for every detector element as a function of θ



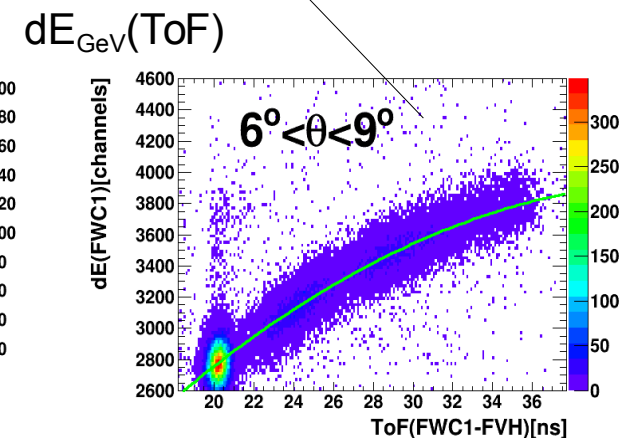
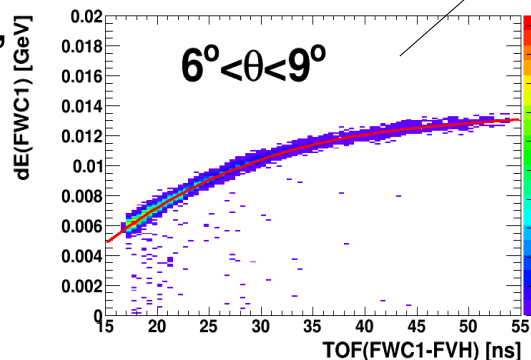
dE Calibration

- Based on ToF
 - **MC:** dE [GeV] vs ToF [ns] $\rightarrow dE_{\text{GeV}}(\text{ToF})$
 - **Data:** dE [channels] vs ToF [ns] $\rightarrow dE_{\text{ch}}(\text{ToF})$
- \rightarrow Run-wise correction, θ -dependency correction

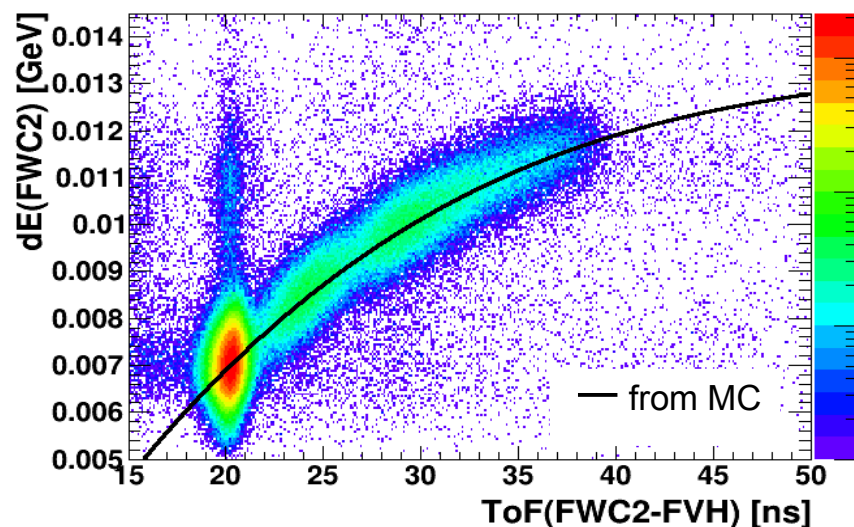


Kinetic Energy Reconstruction

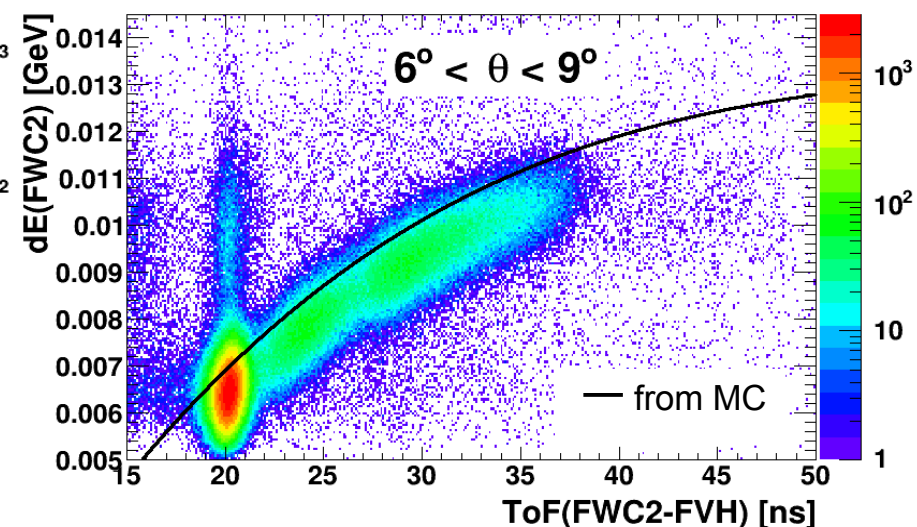
- Based on $E_{\text{kin}}(\text{ToF}_1)$, $E_{\text{kin}}(\text{ToF}_2)$, $E_{\text{kin}}(dE_{\text{FWC1}})$, $E_{\text{kin}}(dE_{\text{FWC2}})$
- χ^2 fit used to obtain the best matching E_{kin}



Energy losses calibration in FWC

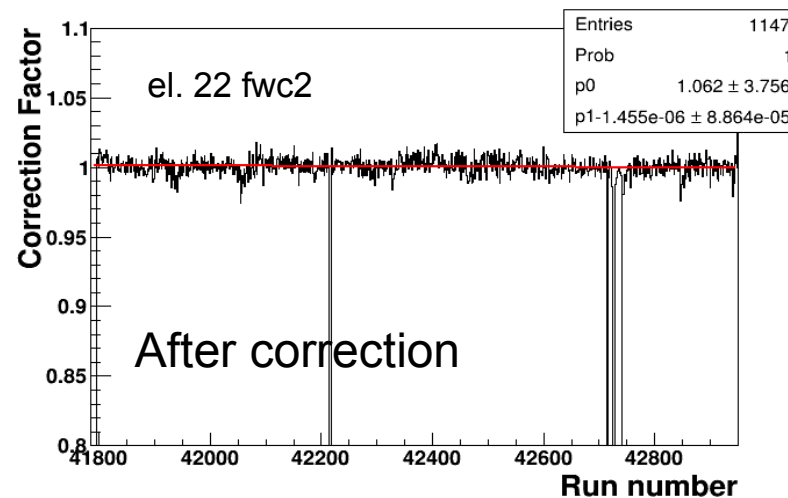
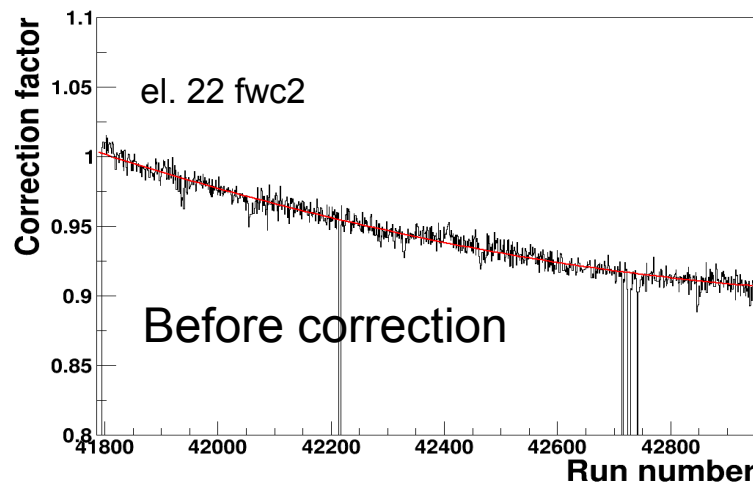


beginning of the beamtime



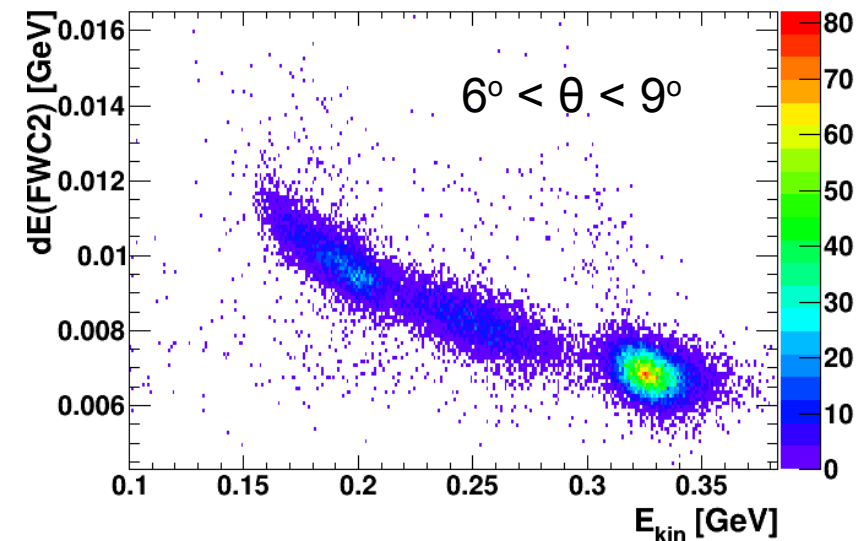
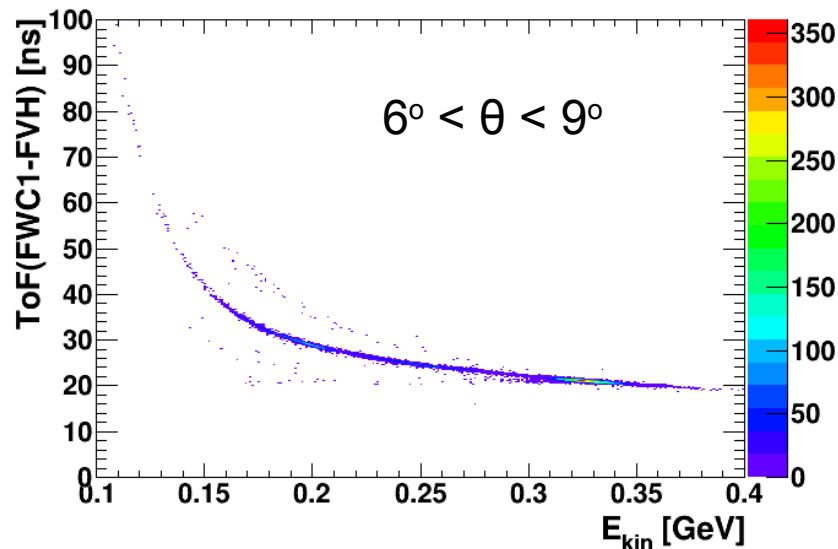
830 runs after beginning of the beamtime
(about $\frac{1}{4}$ of all runs)

- Run correction to dE calibration for every FWC1 and FWC2 element need
- Separate calibration for 2nd part of the beamtime



Kinetic energy calibration

- Minimization of χ^2 :
$$\chi^2 = \sum_{i=1}^n \frac{(dE_i^{meas} - dE(E_{kin})_i)^2}{\sigma_i^2} + \sum_{j=1}^m \frac{(\text{TOF}_j^{meas} - \text{TOF}(E_{kin})_j)^2}{\sigma_j^2}$$
- $E_{kin}(\text{ToF}_1)$, $E_{kin}(\text{ToF}_2)$, $E_{kin}(dE_{FWC1})$, $E_{kin}(dE_{FWC2})$ dependency from MC
- Data based uncertainties of $\text{ToF}(dE)$ as a function of $\text{ToF}(dE)$ (first iteration)

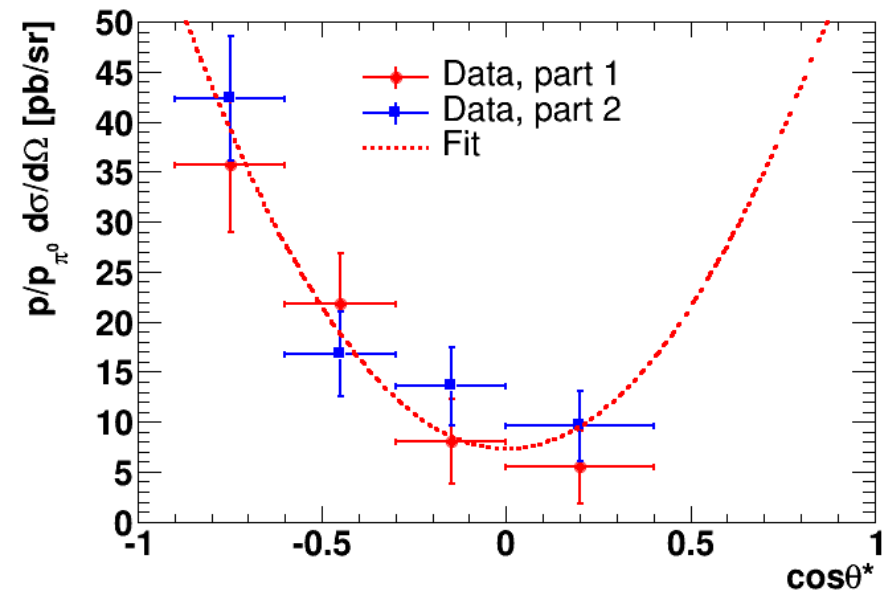
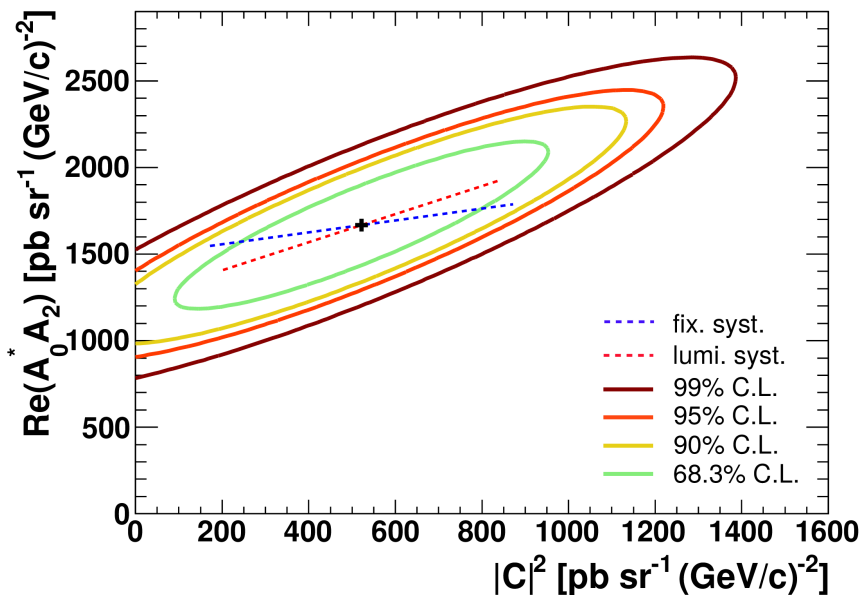


Obtained partial-waves contributions

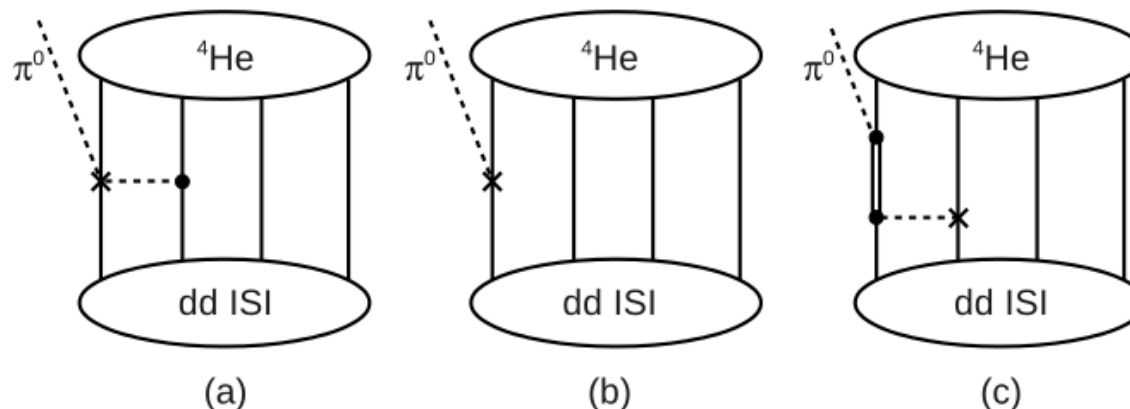
$$\frac{p}{p_{\pi^0}} \frac{d\sigma}{d\Omega} = \left(\underbrace{\frac{2}{3}|A_0|^2_{\text{fixed}}}_{\approx 22 \text{ pb/sr}} - \underbrace{\frac{2}{3}p_{\pi^0}^2 \Re\{A_0^* A_2\}}_{\approx 21 \text{ pb/sr}} + \underbrace{\frac{2}{3}|C|^2 p_{\pi^0}^2}_{\approx 6.6 \text{ pb/sr}} \right) + \left(\underbrace{2p_{\pi^0}^2 \Re\{A_0^* A_2\}}_{\approx 64 \text{ pb/sr}} - \underbrace{\frac{2}{3}|C|^2 p_{\pi^0}^2}_{\approx 6.6 \text{ pb/sr}} \right) \cos^2 \theta^*$$

$$\Re\{A_0^* A_2\} = (1670 \pm 320(\text{stat.})_{-430}^{+80}(\text{syst.})) \text{ pb/} (\text{sr} \cdot (\text{GeV}/c)^2)$$

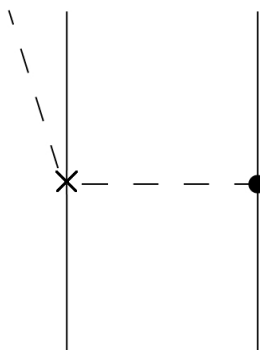
$$|C|^2 = (520 \pm 290(\text{stat.})_{-430}^{+50}(\text{syst.})) \text{ pb/} (\text{sr} \cdot (\text{GeV}/c)^2)$$



Leading diagrams of CSB reactions



Formally leading operators for p -wave pion production in $dd \rightarrow {}^4\text{He}\pi^0$.



- cross – occurrence of CSB
- dot – leading order charge invariant vertex
- dashed line – pions
- single solid line – nucleons
- double solid line – Δ

Leading order diagram for the CSB s-wave amplitudes of the $np \rightarrow d\pi^0$ reaction